
Exercises on Elementary Particle Physics

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1. Representations of $su(2)$

A Lie algebra \mathfrak{g} is a vector space together with a mapping

$$[\cdot, \cdot] : \mathfrak{g} \times \mathfrak{g} \rightarrow \mathfrak{g}$$

satisfying the following conditions:

- (a) The mapping is bilinear (i.e. linear in both entries).
- (b) The mapping is skew-symmetric: $[a, b] = -[b, a]$ for $a, b \in \mathfrak{g}$
- (c) It fulfills the Jacobi identity: $[a, [b, c]] + [b, [c, a]] + [c, [a, b]] = 0$ for $a, b, c \in \mathfrak{g}$

A representation ρ of a Lie algebra \mathfrak{g} on a vector space V is a linear mapping

$$\rho : \mathfrak{g} \rightarrow \text{End}(V)$$

which is an algebra homomorphism, i.e. $\rho([a, b]) = [\rho(a), \rho(b)]$. The dimension of V is called the dimension of the representation: $\dim(\rho) := \dim(V)$.

If there is a vector space $W \subset V$ so that $\rho(W) \subset W$, then the representation is called reducible and W is called the invariant subspace. If such a W does not exist, the representation is called irreducible.

In other words: a representation is irreducible, iff the only invariant subspace is V itself.

As an example, we will concentrate on the algebra $su(2)$ in the following.

- (a) The group $SU(2)$ is the set of all 2-dimensional unitary matrices with determinant 1. Show that the corresponding Lie algebra $su(2)$ is the set of all traceless hermitian matrices. Hint: $\det A = \exp \text{Tr} \log A$.
- (b) Choose the basis

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

for the traceless hermitian matrices. Define

$$J_3 = \frac{1}{2}\sigma_3, \quad J_+ = \frac{1}{2}(\sigma_1 + i\sigma_2), \quad J_- = \frac{1}{2}(\sigma_1 - i\sigma_2),$$

and verify the commutation relations

$$[J_3, J_+] = J_+, \quad [J_3, J_-] = -J_-, \quad [J_+, J_-] = 2J_3.$$

Next, we want to consider **all irreducible, finite-dimensional representations** of $su(2)$ on a vector space V , $\rho(J_i) \in \text{End}(V)$, $i = 3, +, -$. We want to find out, how to classify these representations and which dimensions $\dim(V)$ are allowed.

- (c) Since J_3 is diagonal, $\rho(J_3)$ can also be chosen to be diagonal. Therefore V can be decomposed into eigenspaces of $\rho(J_3)$,

$$V = \bigoplus V_\alpha,$$

where α labels the eigenvalue of $\rho(J_3)$, i.e.

$$(\rho(J_3))v = \alpha v, \quad v \in V_\alpha, \quad \alpha \in \mathbb{C}$$

(shorthand: write J_i for $\rho(J_i)$). Show that $J_+(v) \in V_{\alpha+1}$ and $J_-(v) \in V_{\alpha-1}$

- (d) Prove that all complex eigenvalues α which appear in the above decomposition differ from one another by 1.

Hint: Choose an arbitrary $\alpha_0 \in \mathbb{C}$ from the decomposition and prove that

$$\bigoplus_{k \in \mathbb{Z}} V_{\alpha_0+k} \subset V$$

is indeed equal to V using the irreducibility of the representation.

- (e) Argue that there is $k \in \mathbb{N}$ for which $V_{\alpha_0+k} \neq 0$ and $V_{\alpha_0+k+1} = 0$. Define $n := \alpha_0 + k$. Note that up to now, we only know that $n \in \mathbb{C}$.

Draw a diagram. Write the vector spaces V_{n-2} , V_{n-1} and V_n in a row and indicate the action of J_3 , J_+ and J_- on these vector spaces by arrows.

The eigenvalue n is called highest weight and a vector $v \in V_n$ is called highest weight vector. Is it clear why?

- (f) Choose an arbitrary vector $v \in V_n$ (highest weight vector). Prove that the vectors v, J_-v, J_-^2v, \dots span V .

Hint: Show that the vector space spanned by these vectors is invariant under the action of J_3 , J_+ and J_- and use the irreducibility of the representation.

- (g) Argue that all eigenspaces V_α are 1-dimensional.

(h) Prove that n is a non-negative integer or half-integer and that

$$V = V_{-n} \oplus \dots \oplus V_n .$$

Complement your diagram drawn in part (e). What is the dimension of the representation?

Hint: The representation is finite dimensional, so there exists $m \in \mathbb{N}$ for which $J_-^{m-1}v \neq 0$ and $J_-^m v = 0$. Evaluate the product $J_+ J_-^m v$.

(i) Consider the tensor product of a 2-dimensional and a 3-dimensional irreducible representation of $su(2)$:

$$V = V^{(2)} \otimes V^{(3)}$$

Show that the resulting representation V is reducible and that it can be decomposed into a 2-dim. and a 4-dim. irreducible representation. Shorthand: $\mathbf{2} \otimes \mathbf{3} = \mathbf{2} \oplus \mathbf{4}$.

Hint: The first thing to note is that the action of a Lie algebra on the tensor product of 2 representations is given by: $X(v \otimes w) = Xv \otimes w + v \otimes Xw$, i.e. the eigenvalue of J_3 on V is the sum of the eigenvalues of J_3 on $V^{(2)}$ and $V^{(3)}$. Draw the diagrams of the eigenvalues (with multiplicities). Then use the fact that the eigenspaces of irreducible representations are 1-dimensional.

2. The Lorentz group

The Lorentz group is defined as the set of transformations

$$x^\mu \rightarrow \Lambda^\mu{}_\nu x^\nu$$

which leave the scalar product $\langle x, x \rangle := \eta_{\mu\nu} x^\mu x^\nu$ invariant.

(a) Show that an element λ of the Lie algebra of the Lorentz group satisfies:

$$\lambda^T = -\eta \lambda \eta$$

Hint: Reformulate the statement about the invariance of the scalar product in $\eta_{\mu\nu} = \eta_{\rho\sigma} \Lambda^\rho{}_\mu \Lambda^\sigma{}_\nu$ and write an element of the Lorentz group as $\Lambda^\mu{}_\nu \simeq \delta^\mu{}_\nu - i\lambda^\mu{}_\nu$.

(b) Choose

$$(M^{\mu\nu})^\rho{}_\sigma = i(\eta^{\mu\rho} \delta^\nu{}_\sigma - \eta^{\nu\rho} \delta^\mu{}_\sigma)$$

as a basis for the Lie algebra. What do these matrices look like? Describe the form of the matrices in words. Verify the commutation relations

$$[M^{\mu\nu}, M^{\rho\sigma}] = -i(\eta^{\mu\rho} M^{\nu\sigma} - \eta^{\mu\sigma} M^{\nu\rho} - \eta^{\nu\rho} M^{\mu\sigma} + \eta^{\nu\sigma} M^{\mu\rho}). \quad (1)$$

(c) We split the generators into 2 groups:

$$J^i = \frac{1}{2} \epsilon^{ijk} M^{jk}, \quad K^i = M^{0i}.$$

The J 's have only spatial indices, the K 's have spatial and timelike indices. Verify the commutation relations

$$[J^i, J^j] = i\epsilon^{ijk}J^k, \quad [J^i, K^j] = i\epsilon^{ijk}K^k, \quad [K^i, K^j] = -i\epsilon^{ijk}J^k,$$

and describe the meaning of each relation in words. What kind of transformations do the J 's and K 's correspond to?

- (d) The form of the commutation relations for the Lorentz algebra can still be simplified. Define

$$T_{L/R}^i = \frac{1}{2}(J^i \pm iK^i)$$

and verify the commutation relations

$$[T_L^i, T_L^j] = i\epsilon^{ijk}T_L^k, \quad [T_R^i, T_R^j] = i\epsilon^{ijk}T_R^k, \quad [T_L^i, T_R^j] = 0.$$

- (e) Classify the representations of the Lorentz algebra using what you learned about $su(2)$.

Conclusion: Every representation of the Lorentz algebra can be characterized by 2 non-negative integers or half-integers (j_L, j_R) .

3. Weyl spinors - part I

Summarizing our foregoing considerations, the Lorentz transformation on Minkowski space is given by

$$\Lambda = \exp\left(-\frac{i}{2}\omega_{\mu\nu}M^{\mu\nu}\right). \quad (2)$$

Now we take eq. (1) as the definition of the Lorentz algebra and investigate its representations. To make this point clear, we write $D(\Lambda)$ instead of Λ .

- (a) Define α, β by the equations $\omega_{ij} = \epsilon_{ijk}\alpha_k$ and $\beta_i = \omega_{0i}$ to show

$$D(\Lambda) = \exp\left(-i\left[\vec{\alpha} \cdot \vec{J} + \vec{\beta} \cdot \vec{K}\right]\right) = \exp\left(-i(\vec{\alpha} - i\vec{\beta}) \cdot \vec{T}_L\right) \exp\left(-i(\vec{\alpha} + i\vec{\beta}) \cdot \vec{T}_R\right).$$

Note that T_R^i, T_L^i are still unspecified, we only know their algebra. For a particular representation, we have to make a choice!

- (b) Specialize to a representation: Choose the T_R^i, T_L^i to be the Pauli matrices σ_i . The simplest representations of the Lorentz group are $(1/2, 0)$ and $(0, 1/2)$. An object transforming in the $(1/2, 0)$ representation is called a **left-chiral Weyl spinor**. The definition of a right-chiral Weyl spinor is obvious. How many entries does a Weyl spinor have? Write down the transformation laws for the 2 types of Weyl spinors.

Let D_L, D_R denote the transformation matrices for the left- and right-chiral Weyl spinors.

To be continued ...