
General Relativity

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1. Motion in Schwarzschild solutions

- (a) Write the equations of motion for the Schwarzschild solution.
(Use the Christoffels from exercise sheet 6 with general A, B .)

- (b) Set $\theta = \pi/2$ (why can this be done?), and integrate suitably to get

$$\frac{dt}{dp} = \frac{1}{B(r)}, \quad r^2 \frac{d\phi}{dp} = J = \text{const}. \quad (1)$$

(Multiply the equations for ϕ and t in (a) with $dp/d\phi$ and dp/dt respectively.)

The integral of the equation for r should now give

$$A(r) \left(\frac{dr}{dp} \right)^2 + \frac{J^2}{r^2} - \frac{1}{B(r)} = -E = \text{const}, \quad (2)$$

where p is the parameter along the worldline.

- (c) Show that $d\tau^2 = E dp^2$, and that therefore $E = 0$ must hold for photons, while $E > 0$ for other matter.
- (d) Eliminate dp from the integrals of motion obtained in (b) to get a relation between r and ϕ . Show that

$$\phi = \pm \int \frac{\sqrt{A(r)} dr}{r^2 \sqrt{\frac{1}{J^2 B(r)} - \frac{E}{J^2} - \frac{1}{r^2}}} \quad (3)$$

is a solution.

2. Precession of perihelia

In **1.** we showed that the shape of Schwarzschild trajectories is given by

$$\phi = \pm \int \frac{\sqrt{A(r)} dr}{r^2 \sqrt{\frac{1}{J^2 B(r)} - \frac{E}{J^2} - \frac{1}{r^2}}}, \quad (4)$$

where E and J^2 are constants of the motion. In weak fields we can expand

$$A(r) = 1 + \frac{2MG}{r}, \quad B(r) = 1 - \frac{2MG}{r}.$$

- (a) Determine E and J^2 by looking at the aphelion $r = r_+$ and perihelion $r = r_-$ of a planet in bound orbit around the sun, for general $B(r)$. (At r_{\pm} , $dr/d\phi$ vanishes.)
- (b) Show that the amount of orbital precession per revolution is

$$\Delta\phi = 2|\phi(r_+) - \phi(r_-)| - 2\pi ,$$

where $\phi(r_+) - \phi(r_-)$

$$= \int_{r_-}^{r_+} \left[\frac{r_-^2(B^{-1}(r) - B^{-1}(r_-)) - r_+^2(B^{-1}(r) - B^{-1}(r_+))}{r_+^2 r_-^2 (B^{-1}(r_+) - B^{-1}(r_-))} - \frac{1}{r^2} \right]^{-1/2} \times \frac{\sqrt{A(r)} dr}{r^2} . \quad (5)$$

- (c) Show that for weak fields we can use

$$B^{-1}(r) = 1 + \frac{2MG}{r} + \frac{4M^2G^2}{r^2} ,$$

which makes the first term in Eq. (5) quadratic in $1/r$ and that we can then write

$$\phi(r_+) - \phi(r_-) = \int_{r_-}^{r_+} \left[C \left(\frac{1}{r_-} - \frac{1}{r} \right) \left(\frac{1}{r} - \frac{1}{r_+} \right) \right]^{-1/2} \times \frac{\sqrt{A(r)} dr}{r^2} . \quad (6)$$

- (d) Determine C in the limit $r \rightarrow \infty$. You should get

$$C = 1 - \frac{4MG}{L} + \dots ,$$

where $\frac{1}{L} = \frac{1}{2} \left(\frac{1}{r_+} + \frac{1}{r_-} \right)$. Show that Eq. (5) now reduces to

$$\phi(r_+) - \phi(r_-) = \left(1 + \frac{2MG}{L} \right) \times \int_{r_-}^{r_+} \frac{(1 + \frac{MG}{r}) dr}{r^2 \sqrt{\left(\frac{1}{r_-} - \frac{1}{r} \right) \left(\frac{1}{r} - \frac{1}{r_+} \right)}} . \quad (7)$$

- (e) Calculate $\Delta\phi$. (You can approximate the result of the integral with $(1 + \frac{MG}{L})\pi$.)
- (f) Determine the total precession $\Delta\phi$ for Mercury over the time of a century. (415 revolutions per century; $L = 55.3 \times 10^9$ m; $MG = 1475$ m).
The observed value is (43.11 ± 0.45) arcseconds.