Exercises on Theoretical Particle Physics I Prof. Dr. H.P. Nilles

Due 21.11.2016

8. Yang-Mills theory

(a) Let us take a Lagrangian with a free Dirac fermion

 $\mathscr{L} = \bar{\Psi} i \gamma^{\mu} \partial_{\mu} \Psi$

with Ψ transforming under a global SU(N) as

$$\Psi \to \Psi' = U\Psi$$
, $U = \exp(i\alpha^a T^a)$, $U^{\dagger}U = 1$.

 T^a are the generators of $\mathrm{SU}(N)$ and α^a parametrises the gauge transformation. Show that \mathscr{L} is invariant under this transformation.

 $(1 \ credit)$

(b) As a next step we introduce local SU(N) transformations

 $\Psi \to \Psi' = U(x)\Psi, \quad U(x) = \exp\left(i\,\alpha^a(x)\,T^a\right), \quad U^\dagger(x)U(x) = 1\,.$

Show that the transformation of ${\mathscr L}$ now leads to an extra term

 $\bar{\Psi} U^{\dagger}(x) i \gamma^{\mu} \left(\partial_{\mu} U(x) \right) \Psi$

which means that \mathscr{L} is not invariant under local $\mathrm{SU}(N)$ transformations.

 $(1 \ credit)$

(c) The covariant derivative is defined via the requirement that $D_{\mu}\Psi$ transforms in the same way as Ψ itself

$$D_{\mu}\Psi = \left(\partial_{\mu} + igA_{\mu}^{a}T^{a}\right)\Psi, \qquad D_{\mu}\Psi \to \left(D_{\mu}\Psi\right)' = U(x)D_{\mu}\Psi.$$

Show that this is equivalent to demanding that the gauge boson transforms as

$$A^a_\mu \to \left(A^a_\mu\right)' = A^a_\mu - f^{abc} \alpha^b A^c_\mu - \frac{1}{g} \partial_\mu \alpha^a$$

using a series expansion. How is f^{abc} defined?

 $(1 \ credit)$

(d) Use part (c) to show that the following Lagrangian is gauge invariant

$$\mathscr{L} = \bar{\Psi} i \gamma^{\mu} D_{\mu} \Psi \,.$$

 $(1 \ credit)$

 $(7 \ credits)$

(e) We define the field strength tensor F through

$$ig\left(F^a_{\mu\nu}T^a\right)\Psi = \left(D_{\mu}D_{\nu} - D_{\nu}D_{\mu}\right)\Psi$$

Show that this definition leads to the following expression for its components

$$F^a_{\mu\nu} = \partial_\mu A^a_\nu - \partial_\nu A^a_\mu - g f^{abc} A^b_\mu A^c_\nu.$$
(1 credit)

(f) Use part (e) to derive the transformation property of the field strength tensor

$$F_{\mu\nu} \to (F_{\mu\nu})' = U F_{\mu\nu} U^{-1}$$
$$F^a_{\mu\nu} \to (F^a_{\mu\nu})' = F^a_{\mu\nu} + f^{abc} F^b_{\mu\nu} \alpha^c$$

where $F_{\mu\nu} = F^a_{\mu\nu}T^a$. Because of the last equation the field strength tensor itself is not gauge invariant.

 $(1 \ credit)$

(g) Verify that the product

$$\operatorname{tr}(F_{\mu\nu}F^{\mu\nu})$$

is gauge invariant and thus the gauge invariant Dirac Lagrangian is

$$\mathscr{L} = \bar{\Psi} i \gamma^{\mu} D_{\mu} \Psi - \frac{1}{2} \operatorname{tr} \left(F_{\mu\nu} F^{\mu\nu} \right).$$

9. The Standard Model Higgs effect

(a) The Glashow–Weinberg–Salam theory is the part of the Standard Model (SM) of particle physics which describes the electroweak interactions by a non-Abelian gauge theory with the gauge group $SU(2)_L \times U(1)_Y$. In a one-family approximation, the SM has the following particle content

	$L = \begin{pmatrix} \nu_L \\ e_L \end{pmatrix}$	$R = e_R$	$\Phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}$	$T^a W^a_\mu$	B_{μ}
Hypercharge Y	-1	-2	+1	0	0
$\mathrm{SU}(2)_L$ rep.	2	1	2	3	1
Lorentz rep.	(1/2, 0)	(0, 1/2)	(0,0)	(1/2, 1/2)	(1/2, 1/2)

where L, R contain Dirac spinors and the superscripts in the Higgs doublet denote electromagnetic charges. The corresponding Lagrangian is given by

$$\mathscr{L} = \mathscr{L}_{ ext{gauge}} + \mathscr{L}_{ ext{Higgs}} + \mathscr{L}_{ ext{Yukawa}}$$

with

$$\begin{split} \mathscr{L}_{\text{gauge}} &= \overline{R}(i\gamma^{\mu}D_{\mu})R + \overline{L}(i\gamma^{\mu}D_{\mu})L - \frac{1}{4}F^{a}_{\mu\nu}F^{\mu\nu\,a} - \frac{1}{4}G_{\mu\nu}G^{\mu\nu}\\ \mathscr{L}_{\text{Higgs}} &= (D_{\mu}\Phi)^{\dagger}(D^{\mu}\Phi) - \mu^{2}\Phi^{\dagger}\Phi - \lambda(\Phi^{\dagger}\Phi)^{2}\\ \mathscr{L}_{\text{Yukawa}} &= -G_{e}\left(\overline{L}\Phi R + \overline{R}\Phi^{\dagger}L\right) \end{split}$$

 $(13 \ credits)$

 $(1 \ credit)$

and

$$D_{\mu} = \partial_{\mu} + ig' \frac{Y}{2} B_{\mu} + igT^{a} W^{a}_{\mu}$$

$$G_{\mu\nu} = \partial_{\mu} B_{\nu} - \partial_{\nu} B_{\mu}, \qquad F^{a}_{\mu\nu} = \partial_{\mu} A^{a}_{\nu} - \partial_{\nu} A^{a}_{\mu} + g \epsilon^{abc} A^{b}_{\mu} A^{c}_{\nu}.$$

Write down how the covariant derivative acts on the left- and right-handed lepton doublets/singlet and on the Higgs-doublet.

 $(1 \ credit)$

(b) Show that the Lagrangian \mathscr{L} given in part (a) is Lorentz invariant.

 $(2 \ credits)$

(c) Show that \mathscr{L} is also gauge invariant.

 $(2 \ credits)$

(d) For the Higgs mechanism to work we need $\mu^2 < 0$. For which value of $|\Phi|$ does the Higgs potential obtain a minimum? In unitary gauge we can choose

$$\Phi = \frac{1}{\sqrt{2}} \begin{pmatrix} 0\\ v + \eta(x) \end{pmatrix}$$

where v is the vacuum expectation value and $\eta(x)$ a real field. Show that the Higgs potential around the minimum in the unitary gauge is given by

$$V(\Phi) = -\mu^2 \eta^2(x) + \lambda v \, \eta^3(x) + \frac{\lambda}{4} \eta^4(x) \,.$$

What is the mass of the η field? Compare the degrees of freedom in the Higgs sector before and after spontaneous symmetry breaking.

 $(2 \ credits)$

(e) Consider the kinetic energy terms of the Higgs field in $\mathscr{L}_{\text{Higgs}}$. Show that

$$\begin{split} (D_{\mu}\Phi)^{\dagger}(D^{\mu}\Phi) &= \frac{1}{2}\partial_{\mu}\eta\,\partial^{\mu}\eta + \frac{1}{4}g^{2}\left(v+\eta\right)^{2}W_{\mu}^{-}W^{+\,\mu} \\ &+ \frac{1}{8}\left(v+\eta\right)^{2}\left(W_{\mu}^{3}\ B_{\mu}\right) \begin{pmatrix} g^{2} & -g'g \\ -g'g & g'^{2} \end{pmatrix} \begin{pmatrix} W^{3\,\mu} \\ B^{\mu} \end{pmatrix} \\ \text{with } W^{\pm\,\mu} &= \frac{1}{\sqrt{2}}(W^{1\,\mu}\mp i\,W^{2\,\mu}). \end{split}$$

 $(2 \ credits)$

(f) What are the masses for W^{\pm}_{μ} ? To see the masses of W^{3}_{μ} and B_{μ} bosons one has to diagonalize the matrix given in part (e)

$$\frac{v^2}{8} \begin{pmatrix} W^3_{\mu} & B_{\mu} \end{pmatrix} \mathcal{O}^T \mathcal{O} \begin{pmatrix} g^2 & -g'g \\ -g'g & g'^2 \end{pmatrix} \mathcal{O}^T \mathcal{O} \begin{pmatrix} W^{3\mu} \\ B^{\mu} \end{pmatrix} = \begin{pmatrix} Z_{\mu} & A_{\mu} \end{pmatrix} \begin{pmatrix} m_Z^2 & 0 \\ 0 & m_A^2 \end{pmatrix} \begin{pmatrix} Z^{\mu} \\ A^{\mu} \end{pmatrix}$$

Determine this orthogonal matrix \mathcal{O} by computing the corresponding eigenvalues and eigenvectors. What are the masses of the Z_{μ} and A_{μ} fields? Compare the degrees of freedom in the gauge sector to the situation before the symmetry breakdown. What can you say about the total amount of degrees of freedom?

 $(2 \ credits)$

(g) As you know, an orthogonal 2×2 matrix can be written as

$$\mathcal{O} = \begin{pmatrix} \cos \theta_W & -\sin \theta_W \\ \sin \theta_W & \cos \theta_W \end{pmatrix}.$$

Write $\cos \theta_W$ in terms of g' and g. Show for the ratio of the W- and Z-boson masses

$$\frac{m_W}{m_Z} = \cos \theta_W$$

The angle θ_W is called Weinberg angle or weak mixing angle.

 $(1 \ credit)$

(h) Finally, we want to consider the covariant derivative again. Substitute the fields B_{μ} and W^{a}_{μ} by W^{\pm}_{μ} , Z_{μ} and A_{μ} and show

$$D_{\mu} = \partial_{\mu} + i A_{\mu} e Q + i Z_{\mu} \frac{1}{\sqrt{g^2 + g'^2}} \left(g^2 T_3 - g'^2 \frac{Y}{2}\right) + \frac{ig}{\sqrt{2}} \begin{pmatrix} 0 & W_{\mu}^+ \\ W_{\mu}^- & 0 \end{pmatrix}$$

where we have defined the electric charge $e = \frac{g'g}{\sqrt{g^2 + g'^2}}$ and $Q = T_3 + \frac{Y}{2}$.

 $(1 \ credit)$