## Advanced Quantum Theory (WS 24/25) Homework no. 10 (December 9, 2024) Please hand in your solution by Sunday, December 15!

## Quickies

Q1: (i) Does the Born approximation for the scattering cross section depend on the sign of the scattering potential (attractive or repulsive)? (ii) Does the existence of bound states depend on this sign? (iii) What does this imply for the ability of the Born approximation to predict resonances in the scattering cross section? [3P]

**Q2:** At small momenta, the total cross section for the  $\ell$ th partial wave is  $\propto |\vec{k}|^{4\ell}$ . What does this imply for the scaling of (i) the partial wave amplitude  $a_{\ell}$ ; (ii) the scaling of the scattering phase  $\delta_{\ell}$  with  $|\vec{k}|$  at small  $|\vec{k}|$ ? [2P]

Q3: Can the scattering formalism developed in class so far (scattering of a non-relativistic particle on a fixed scattering center) be applied to the following reactions, and, if so, for what range of energy of the incoming particle: (i) electron-proton scattering; (ii) neutron scattering on a lead nucleus; (iii) proton-proton scattering? [3P]

## 1) Potential Well and Born Approximation

In this exercise we analyze scattering on a spherically symmetric potential well, as in sec. 5.5 in class, i.e.  $V(r) = -V_0$  for  $r \leq r_0$  while V(r) = 0 for  $r > r_0$ . M is the mass of the scattering particle.

1. Show that in Born approximation the scattering amplitude  $f_{\vec{k}}(\theta,\phi)$  can be written as

$$f_{\vec{k}}^{\text{Born}} = \frac{2MV_0r_0}{\hbar^2 |\vec{q}|^2} \left[ \frac{1}{r_0 |\vec{q}|} \sin(r_0 |\vec{q}|) - \cos(r_0 |\vec{q}|) \right], \tag{1}$$

where  $\vec{q} = \vec{k}' - \vec{k}$ , i.e.  $\hbar \vec{q}$  is the momentum exchanged in the scattering. Hint: Recall eq.(5.22) in class,

$$f_{\vec{k}}^{\text{Born}} = -\frac{2M}{\hbar^2 |\vec{q}|} \int_0^\infty r V(r) \sin(|\vec{q}|r) dr \,. \tag{2}$$

[3P]

2. Show that for  $|\vec{q}|r_0 \ll 1$ ,

 $f_{\vec{k}}^{\text{Born}} \simeq v_0 \frac{r_0}{3} \,, \tag{3}$ 

where

$$v_0 = \frac{2MV_0r_0^2}{\hbar^2}\,,\tag{4}$$

see eq.(5.58) in class. *Hint:* You need to expand the trigonometric functions in eq.(1) up to the second term. [3P]

3. Show that therefore in the same limit, the total scattering cross section is

$$\sigma^{\text{Born}} \simeq 4\pi r_0^2 \frac{v_0^2}{9} \,. \tag{5}$$

Interpret this result. In particular, show that in this limit only the S-wave contributes, and that  $|\vec{q}|r_0 \ll 1$  requires  $|\vec{k}|r_0 \ll 1$ . [3P]

4. Recall that the Born approximation means that one uses the incoming (plane) wave when computing the contribution of the spherical wave to the entire wave function; from eq(5.12), this contribution is given in Born approximation by

$$\psi_{\rm scat}^{\rm Born}(\vec{x}) = -\frac{2M}{\hbar} \int d^3x' \frac{{\rm e}^{i|\vec{k}||\vec{x}-\vec{x}'|}}{4\pi |\vec{x}-\vec{x}'|} V(\vec{x}') {\rm e}^{i\vec{k}\cdot\vec{x}'} \,. \tag{6}$$

Argue that the Born approximation can only be trusted if  $|\psi_{\text{scat}}^{\text{Born}}(\vec{x}=0)| \ll 1$ . Apply this to our potential well in the limit  $|\vec{k}|r_0 \ll 1$ , and show that it is equivalent to

$$\frac{v_0}{2} \ll 1 \,, \tag{7}$$

where  $v_0$  has been defined in eq.(4).

- [5P]
- 5. For the case at hand the scattering cross section can be computed exactly. In particular, from eq.(5.69) in class we have for the S-wave scattering phase:

$$\delta_0 = \arctan\left[\frac{k}{q}\tan(qr_0)\right] - kr_0, \qquad (8)$$

where

$$qr_0 = \sqrt{(kr_0)^2 + v_0}, \qquad (9)$$

see eq.(5.70) in class, i.e.  $\hbar q$  is *not* the absolute value of the momentum exchange. Show that eq.(8) reproduces eq.(5) in the appropriate limit  $kr_0$ ,  $v_0 \ll 1$ . *Hint:* You again need to keep the two leading terms in the expansion of the trigonometric functions in eq.(8). [5P]